## CS 258: Quantum Cryptography

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Previously...

## Short Integer Solution (SIS)

Input:  $\mathbf{A} \leftarrow \mathbb{Z}_q^{n \times m}$  (short, wide)

Chosen uniformly at random

**Goal:** find vector  $\mathbf{x} \in \mathbb{Z}^m$  such that:

$$\mathbf{A} \cdot \mathbf{x} \mod q = 0$$

$$0 < |\mathbf{x}| \le \beta$$

#### Search LWE

Input: 
$$\mathbf{A} \leftarrow \mathbb{Z}_q^{n imes m}$$
 (short, wide) Chosen uniformly at random  $\mathbf{u} = \mathbf{A}^T \cdot \mathbf{s} + \mathbf{e} \bmod q$  where  $\mathbf{s}$  uniform in  $\mathbb{Z}_q^n$   $\mathbf{e} \leftarrow D_\sigma^m$ 

Output: s (in this regime, s is whp unique)

**Thm** (restated): If SIS cannot be solved in quantum polynomial time for  $\beta=mq/2\sigma$ , then neither can decision LWE with error  $\sigma$ 

Now used to justify hardness of LWE

Even earlier...

#### **Group Action**

An (abelian) group action is a triple  $(\mathbb{G},\mathcal{X},*)$  where:

- G is an (abelian) group, written additively
- $oldsymbol{\cdot} \mathcal{X}$  is a set
- $ullet *: \mathbb{G} imes \mathcal{X} o \mathcal{X}$  is an efficient binary operation satisfying

$$g * (h * x) = (g+h) * x$$

- There is some element  $x_0 \in \mathcal{X}$  that can be efficiently computed
- Usually ask that for each  $x,y\in\mathcal{X}$  , there exists a unique  $g\in\mathbb{G}$  such that y=g\*x
- Also usually ask that it is possible to efficiently identify elements of  ${\mathcal X}$

**Thm** [Kuperberg]: Dlog in (abelian) group actions can be solved in time  $2^{O(\sqrt{\log q})}$ , where q is the group order

#### Broader Picture: Hidden Shifts

Kuperberg actually solves a much more general problem called hidden shift

Given 
$$f_0, f_1: \mathbb{G} o \mathcal{X}$$
 injective, such that  $f_1(g) = f_0(a+g)$  , find  $a$  ( $\mathbb{G}$  written additively)

Group action Dlog is a special case of hidden shift where

$$f_0(g) = g * x_0$$
  $f_1(g) = g * x_1 = (g + a) * x_0$ 

Today: More Quantum Algorithms for Lattices

#### LWE as Hidden Shift

Suppose for the moment that LWE had no error

Input: 
$$\mathbf{A} \leftarrow \mathbb{Z}_q^{n \times m}$$
 (short, wide) Chosen uniformly at random  $\mathbf{u} = \mathbf{A}^T \cdot \mathbf{s} \bmod q$  where  $\mathbf{s}$  uniform in  $\mathbb{Z}_q^n$ 

Output: s (in this regime, s is whp unique)

Of course, this is easy due by Gaussian elimination

#### LWE as Hidden Shift

$$f_0(\mathbf{r}) = \mathbf{A}^T \cdot \mathbf{r} \bmod q$$

$$f_1(\mathbf{r}) = \mathbf{A}^T \cdot \mathbf{r} + \mathbf{u} \mod q = \mathbf{A}^T \cdot (\mathbf{r} + \mathbf{s}) \mod q = f_0(\mathbf{r} + \mathbf{s} \mod q)$$

So solving hidden shift allows us to recover s

#### Ok, but what about the error e?

Solution: round

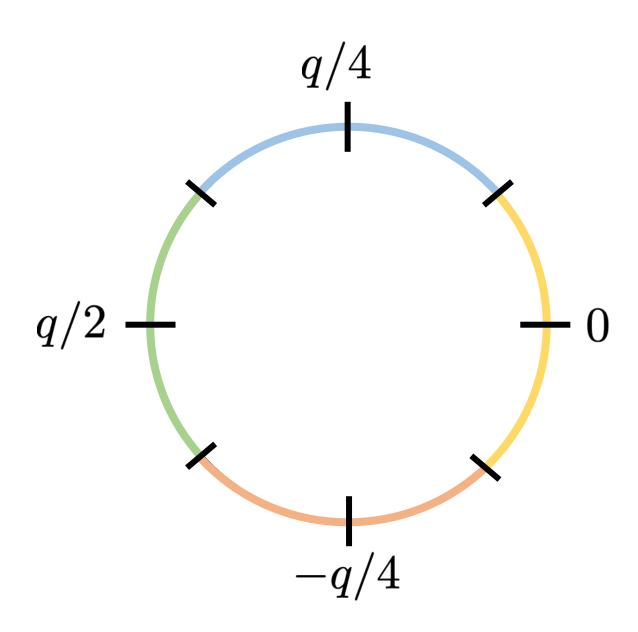
Output closest of 
$$-q/4,0,q/4,q/2$$

$$f_0(\mathbf{r}) = \lfloor \mathbf{A}^T \cdot \mathbf{r} \bmod q \rceil_{q/4}$$

$$f_1(\mathbf{r}) = \lfloor \mathbf{A}^T \cdot \mathbf{r} + \mathbf{u} \bmod q \rceil_{q/4}$$

Idea: if error small enough, rounding eliminates error

$$\lfloor x + e \rceil_{q/4} = \lfloor x \rceil_{q/4}$$
 typically if  $e$  small



Now if  $\mathbf{u} = \mathbf{A}^T \cdot \mathbf{s} + \mathbf{e} \bmod q$ 

$$f_1(\mathbf{r} = \lfloor \mathbf{A}^T \cdot \mathbf{r} + \mathbf{u} \bmod q \rceil_{q/4})$$

$$= \lfloor \mathbf{A}^T \cdot (\mathbf{r} + \mathbf{s}) + \mathbf{e} \bmod q \rceil_{q/4}$$

$$= ? \lfloor \mathbf{A}^T \cdot (\mathbf{r} + \mathbf{s}) \bmod q \rceil_{q/4}$$

$$= f_0(\mathbf{r} + \mathbf{s} \bmod q)$$

#### Need to show:

- Rounding actually gets rid of  ${f e}$   $f_0, f_1$  are injective

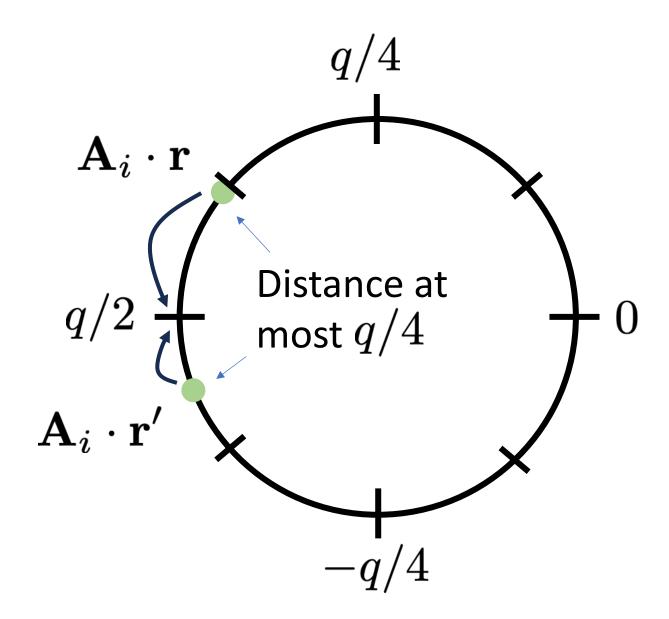
#### Injectivity

Suffices to only look at  $f_0$  , as hidden shift property will imply injectivity for  $f_1$ 

$$f_0(\mathbf{r}) = f_0(\mathbf{r}') \iff [\mathbf{A}^T \cdot \mathbf{r} \bmod q]_{q/4} = [\mathbf{A}^T \cdot \mathbf{r}' \bmod q]_{q/4}$$
$$\implies |\mathbf{A}^T \cdot (\mathbf{r} - \mathbf{r}')|_{\infty} \le q/4$$

Max of absolute values of entries

## Injectivity



#### Injectivity

$$f_0({f r})=f_0({f r}')$$
 ,  ${f r}
eq {f r}'$ 

$$\exists \mathbf{v} : |\mathbf{A}^T \cdot \mathbf{v} \bmod q|_{\infty} \le q/4$$

Claim: with overwhelming probability over  ${f A}$  , no such  ${f v}$ 

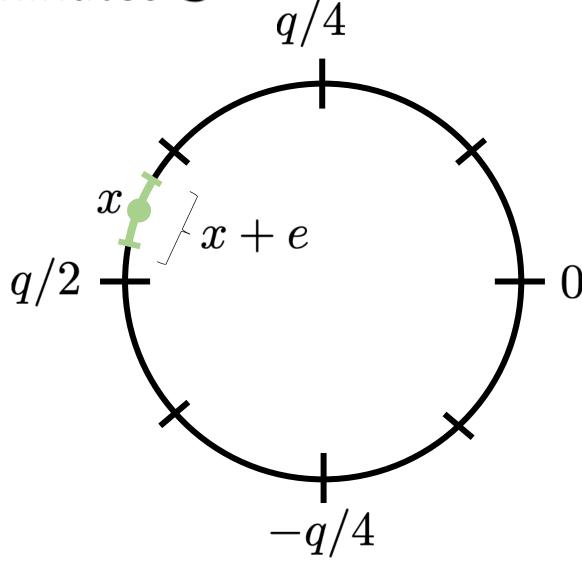
Proof: for any 
$$\mathbf{v}$$
,  $\Pr[|\mathbf{A}_i \cdot \mathbf{v} \bmod q| \le q/4] = 1/2$ 

$$\longrightarrow \Pr[|\mathbf{A} \cdot \mathbf{v} \bmod q|_{\infty} \le q/4] = 2^{-m} = 2^{-\Omega(n \log q)}$$

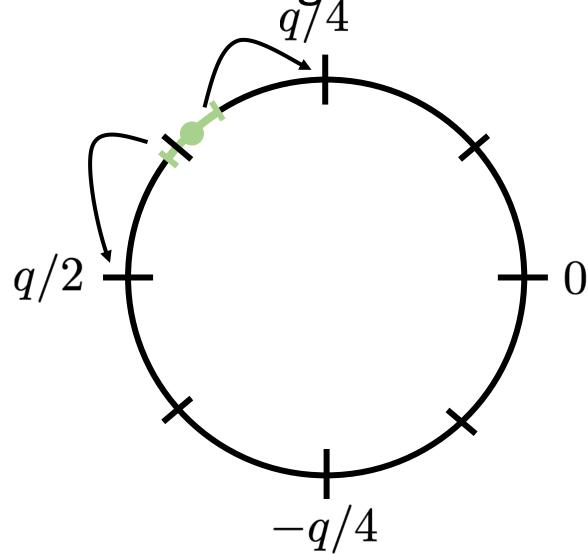
Union-bound over all  $2^{n \log q}$  choices of  ${f v}$ 

$$ightharpoonup \Pr[\exists \mathbf{v} : |\mathbf{A}^T \cdot \mathbf{v} \bmod q|_{\infty} \le q/4] \le 2^{-\Omega(n \log q)}$$

Rounding eliminates **e** 



The problem with rounding



## The problem with rounding

Each entry has a  $\approx O(\sigma/q)$  chance of being too close to a rounding boundary

Over m entries, probability of some error is  $pprox O(\sigma m/q)$ 

#### Can we apply Kuperberg?

• Prepare 
$$\frac{1}{\sqrt{2q^n}}\sum_{\mathbf{r}\in\mathbb{Z}_q^n,b\in\{0,1\}}|\mathbf{r},b\rangle_{\mathcal{A}}|0\rangle_{\mathcal{B}}$$

• Apply 
$$U_f$$
 where  $f(\mathbf{r},b)=f_b(\mathbf{r})$ : 
$$\frac{1}{\sqrt{2q^n}}\sum_{\mathbf{r}\in\mathbb{Z}_q^n,b\in\{0,1\}}|\mathbf{r},b\rangle_{\mathcal{A}}|f_b(\mathbf{r})\rangle_{\mathcal{B}}$$

$$= \frac{1}{\sqrt{2q^n}} \sum_{\mathbf{r} \in \mathbb{Z}_q^n, b \in \{0,1\}} |\mathbf{r}, b\rangle_{\mathcal{A}} ||\mathbf{A}^T \cdot (\mathbf{r} + b\mathbf{s}) + b\mathbf{e} \bmod q|_{q/4}\rangle_{\mathcal{B}}$$

#### Can we apply Kuperberg?

$$\frac{1}{\sqrt{2q^n}} \sum_{\mathbf{r} \in \mathbb{Z}_q^n, b \in \{0,1\}} |\mathbf{r}, b\rangle_{\mathcal{A}} ||\mathbf{A}^T \cdot (\mathbf{r} + b\mathbf{s}) + b\mathbf{e} \bmod q|_{q/4}\rangle_{\mathcal{B}}$$

• Measure  $\mathcal{B} \rightarrow$  Measurement outcome zState collapses to  ${f r}, b$  consistent with z

If  $\mathbf{A}^T \cdot (\mathbf{r} + b\mathbf{s}) \bmod q$  is far from rounding boundary,  $[\mathbf{A}^T \cdot (\mathbf{r} + b\mathbf{s}) + b\mathbf{e} \bmod q]_{q/4} = [\mathbf{A}^T \cdot (\mathbf{r} + b\mathbf{s}) \bmod q]_{q/4}$ 



State collapses to  $\frac{1}{\sqrt{2}}|\mathbf{r},0\rangle + \frac{1}{\sqrt{2}}|\mathbf{r}-\mathbf{s} \bmod q,1\rangle$ 



#### Possible issues with applying Kuperberg

1. The shift lives in  $\mathbb{Z}_q^n$  instead of  $\mathbb{Z}_{2^n}$ 

Turns out to not be a problem

2. The errors

Big problem!!!

If  $\mathbf{A}^T \cdot (\mathbf{r} + b\mathbf{s}) \bmod q$  is **close** to rounding boundary,

$$\lfloor \mathbf{A}^T \cdot (\mathbf{r} + b\mathbf{s}) + b\mathbf{e} \bmod q \rceil_{q/4} \neq \lfloor \mathbf{A}^T \cdot (\mathbf{r} + b\mathbf{s}) \bmod q \rceil_{q/4}$$



Recall next step of Kuperberg: apply  $\mathsf{QFT}_q$  to first register, measure

$$\frac{1}{\sqrt{q^n}} \sum_{\mathbf{t}} |\mathbf{t}, b\rangle e^{i2\pi \mathbf{r} \cdot \mathbf{t}/q} \implies |\mathbf{t}, b\rangle$$

#### **Combining Samples**

Bad sample

$$\mathbf{t}_{0}, |\psi_{0}\rangle = \frac{1}{\sqrt{2}}|0\rangle + \frac{1}{\sqrt{2}}e^{-i2\pi\mathbf{s}\cdot\mathbf{t}_{0}/q}|1\rangle \qquad \mathbf{t}_{1}, |\psi_{1}\rangle = |b\rangle$$

$$|\mathbf{t}_1,|\psi_1
angle=|b
angle$$

$$\mathsf{CNOT}|\psi_0\rangle|\psi_1\rangle = \frac{1}{\sqrt{2}}|0,b\rangle + \frac{1}{\sqrt{2}}e^{-i2\pi\mathbf{s}\cdot\mathbf{t}_0/q}|1,1-b\rangle$$

Measure second qubit:  $|0\rangle$  or  $|1\rangle$ 

Combining with bad samples gives bad samples

Kuperberg requires  $2^{O(\sqrt{\log(q^n)})} = 2^{O(\sqrt{n\log q})}$  samples

If any of those samples are bad, Kuperberg fails

ightharpoonup Need  $\sigma m/q=2^{-\Omega(\sqrt{n\log q})}$  to have decent chance of all samples being good

It turns out that, in this regime, classical attacks already exist

Significant open question: can Kuperberg's algorithm be made robust to errors?

A positive solution would give a subexponential-time attack on LWE, which would give lattice crypto a significant efficiency penalty

Even beyond LWE, making robust to errors could be important for realizing Kuperberg on a realistic quantum computer

## Other possible algorithms

Quasi-polynomial attack on hidden shifts over  $\mathbb{Z}_q^n$  , when  $q=2^r$ 

Note that for LWE, hardness is robust to modulus, and can take it to be power of 2

$$\mathbf{t}_{j}, |\psi_{\mathbf{t}_{j}}\rangle = \frac{1}{\sqrt{2}}|0\rangle + \frac{1}{\sqrt{2}}e^{-i2\pi\mathbf{s}\cdot\mathbf{t}_{j}/q}|1\rangle$$

Write 
$$|\psi_{\mathbf{t}_1}\rangle|\psi_{\mathbf{t}_2}\rangle\cdots$$
 as

$$\frac{1}{\sqrt{2^{\ell}}} \sum_{\mathbf{b} \in \{0,1\}^{\ell}} |\mathbf{b}\rangle e^{-i2\pi \mathbf{s}^T \mathbf{T} \mathbf{b}/q}$$

Where 
$$\mathbf{T} = (\mathbf{t}_1 \ \mathbf{t}_2 \ \cdots \ \mathbf{t}_\ell)$$

$$\frac{1}{\sqrt{2^{\ell}}} \sum_{\mathbf{b} \in \{0,1\}^{\ell}} |\mathbf{b}\rangle e^{-i2\pi \mathbf{s}^T \mathbf{T} \mathbf{b}/q}$$

Let's assume mod 2 that  ${f T}$  has a 1-dimensional kernel Will be true if we choose  $\ell pprox n+1$ 

$$\frac{1}{\sqrt{2^{\ell}}} \sum_{\mathbf{b} \in \{0,1\}^{\ell}} |\mathbf{b}\rangle e^{-i2\pi \mathbf{s}^T \mathbf{T} \mathbf{b}/q}$$

Now apply map  $|\mathbf{b}\rangle\mapsto |\mathbf{b},\mathbf{Tb} \bmod 2\rangle$  , and measure second register  $\to \mathbf{z}$ 

$$\frac{1}{\sqrt{2}}|\mathbf{b}_{0}\rangle e^{-i2\pi\mathbf{s}^{T}\mathbf{T}\mathbf{b}_{0}/q} + \frac{1}{\sqrt{2}}|\mathbf{b}_{1}\rangle e^{-i2\pi\mathbf{s}^{T}\mathbf{T}\mathbf{b}_{1}/q}$$

$$= e^{-i2\pi\mathbf{s}^{T}\mathbf{T}\mathbf{b}_{0}/q} \left(\frac{1}{\sqrt{2}}|\mathbf{b}_{0}\rangle + \frac{1}{\sqrt{2}}|\mathbf{b}_{1}\rangle e^{-i2\pi\mathbf{s}^{T}\mathbf{T}(\mathbf{b}_{1}-\mathbf{b}_{0})/q}\right)$$

Where  $\mathbf{b}_0, \mathbf{b}_1$  are the two values with

$$\mathbf{Tb}_0 \mod 2 = \mathbf{Tb}_1 \mod 2 = \mathbf{z}$$

$$e^{-i2\pi\mathbf{s}^T\mathbf{T}\mathbf{b}_0/q}\left(\frac{1}{\sqrt{2}}|\mathbf{b}_0\rangle + \frac{1}{\sqrt{2}}|\mathbf{b}_1\rangle e^{-i2\pi\mathbf{s}^T\mathbf{T}(\mathbf{b}_1-\mathbf{b}_0)/q}\right)$$

Now map 
$$|\mathbf{b}_0\rangle \mapsto |0\rangle, |\mathbf{b}_1\rangle \mapsto |1\rangle$$

$$e^{-i2\pi\mathbf{s}^T\mathbf{T}\mathbf{b}_0/q}\left(\frac{1}{\sqrt{2}}|0\rangle + \frac{1}{\sqrt{2}}|1\rangle e^{-i2\pi\mathbf{s}^T\mathbf{T}(\mathbf{b}_1-\mathbf{b}_0)/q}\right)$$

Global phase doesn't matter:  $|\psi_{\mathbf{T}(\mathbf{b}_1 - \mathbf{b}_0)}\rangle$ 

Now, observe that  ${f T}({f b}_1-{f b}_0)$  is even, say  $2{f t}'$ 

$$\frac{1}{\sqrt{2}}|0\rangle + \frac{1}{\sqrt{2}}|1\rangle e^{-i2\pi\mathbf{s}\cdot2\mathbf{t}'/q} = \frac{1}{\sqrt{2}}|0\rangle + \frac{1}{\sqrt{2}}|1\rangle e^{-i2\pi\mathbf{s}\cdot\mathbf{t}'/(q/2)}$$

Reduced the modulus by factor of 2

## Each step divides number of samples by pprox n divides modulus by 2

Number of samples needed:

$$\approx n^{\log q} = 2^{(\log n)(\log q)}$$

For LWE parameters, this is  $2^{O(\log^2 n)}$  , quasi-polynomial!

But, errors still break this algorithm

## Multiple shifts

#### Multiple shifts

$$f_0$$
  
 $f_1(\mathbf{r}) = f_0(\mathbf{r} + \mathbf{s})$   
 $f_2(\mathbf{r}) = f_0(\mathbf{r} + 2\mathbf{s})$ 

If we could go all the way to  $f_q$  , we'd actually get a periodic function. Maybe something in between makes the problem easier?

• • •

### Multiple shifts for LWE

$$f_j(\mathbf{r}) = \lfloor \mathbf{A}^T \cdot \mathbf{r} + j\mathbf{u} \bmod q \rceil_{q/4} = \lfloor \mathbf{A}^T \cdot (\mathbf{r} + j\mathbf{s}) + j\mathbf{e} \bmod q \rceil_{q/4}$$

Larger j means larger errors → definitively can't get all the way to periodic

To date, no attack on LWE based on any of these ideas

# Next time: when using post-quantum building blocks is not enough